



# UNB application of Stokes-Helmert's approach to geoid computation

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# Outline



- Formulation of the appropriate BVP
- Determination of the boundary values
- Rigorous treatment of topo-effects + DWC
- Modified Stokes's formula
- Practical results & problems
- Current developments
- Conclusions

**Unknown function**

**Region of interest**

$$T(r, \Omega) = W(r, \Omega) - U(r, \Omega), \quad \forall r \geq r_g, \Omega \in \Omega_0.$$

**Ideal world**

$$N(\Omega) = \frac{T(r_g, \Omega)}{\gamma_0(\phi)}$$

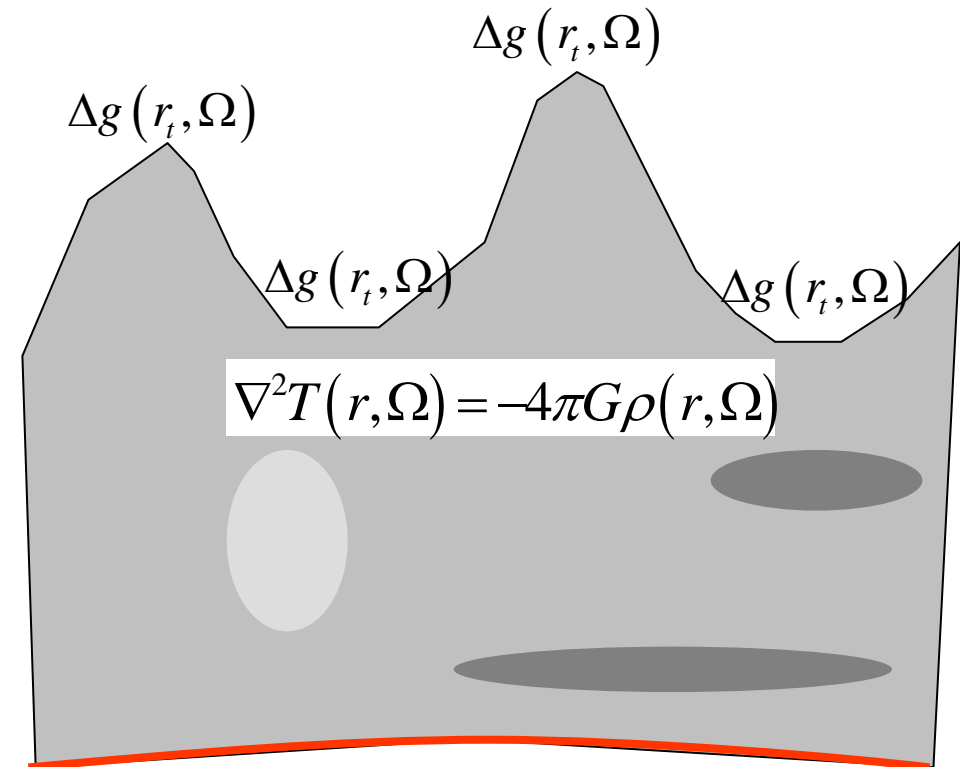
**Real world**

$$\lim_{r \rightarrow \infty} T(r, \Omega) = O(r^{-3})$$

$$\nabla^2 T(r, \Omega) = 0$$

$$\Delta g(r_g, \Omega) \quad \Delta g(r_g, \Omega) \quad \Delta g(r_g, \Omega)$$

**Geoid**





# UNB Helmertization of the disturbing potential



$$T^h(r, \Omega) = T(r, \Omega) - V^t(r, \Omega) + V^{ct}(r, \Omega) - V^a(r, \Omega) + V^{ca}(r, \Omega)$$

$$\nabla^2 T^h(r, \Omega) = 0$$

## Fundamental gravimetric equation

$$\Delta g(r, \Omega) = -\frac{\partial T(r, \Omega)}{\partial n} + \gamma[r - \zeta(r, \Omega), \phi]^{-1} \frac{\partial \gamma(r, \phi)}{\partial n} T(r, \Omega)$$

## Conversion of free-air to Helmert's anomaly

$$\begin{aligned} \Delta g^h(r, \Omega) = \Delta g^{FA}(r, \Omega) &+ \frac{\partial [V^t(r, \Omega) - V^{ct}(r, \Omega)]}{\partial r} + \frac{\partial [V^a(r, \Omega) - V^{ca}(r, \Omega)]}{\partial r} + \\ &+ \frac{2}{r} [V^t(r, \Omega) - V^{ct}(r, \Omega)] + \frac{2}{r} [V^a(r, \Omega) - V^{ca}(r, \Omega)] + \varepsilon_{\delta g} + \varepsilon_n \end{aligned}$$

Harmonic!!!  $\nabla^2 [r \Delta g^h(r, \Omega)] = 0$



# UNB “Spherical” topographic effects



Planar

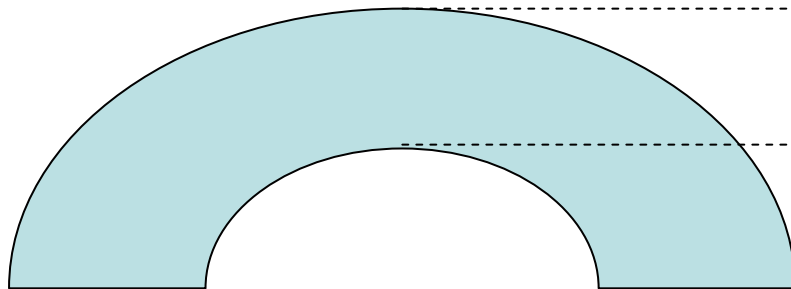


$$V^t(r_t, \Omega) \rightarrow \infty \quad A(r_t, \Omega) = 2\pi G\rho H$$

$$V^t(r_g, \Omega) \rightarrow \infty \quad A(r_g, \Omega) = -2\pi G\rho H$$

Disqualifies for the harmonization of the Earth gravity field!!!

Spherical



$$V^t(r_g, \Omega) = 4\pi G\rho RH + O(H/R)$$

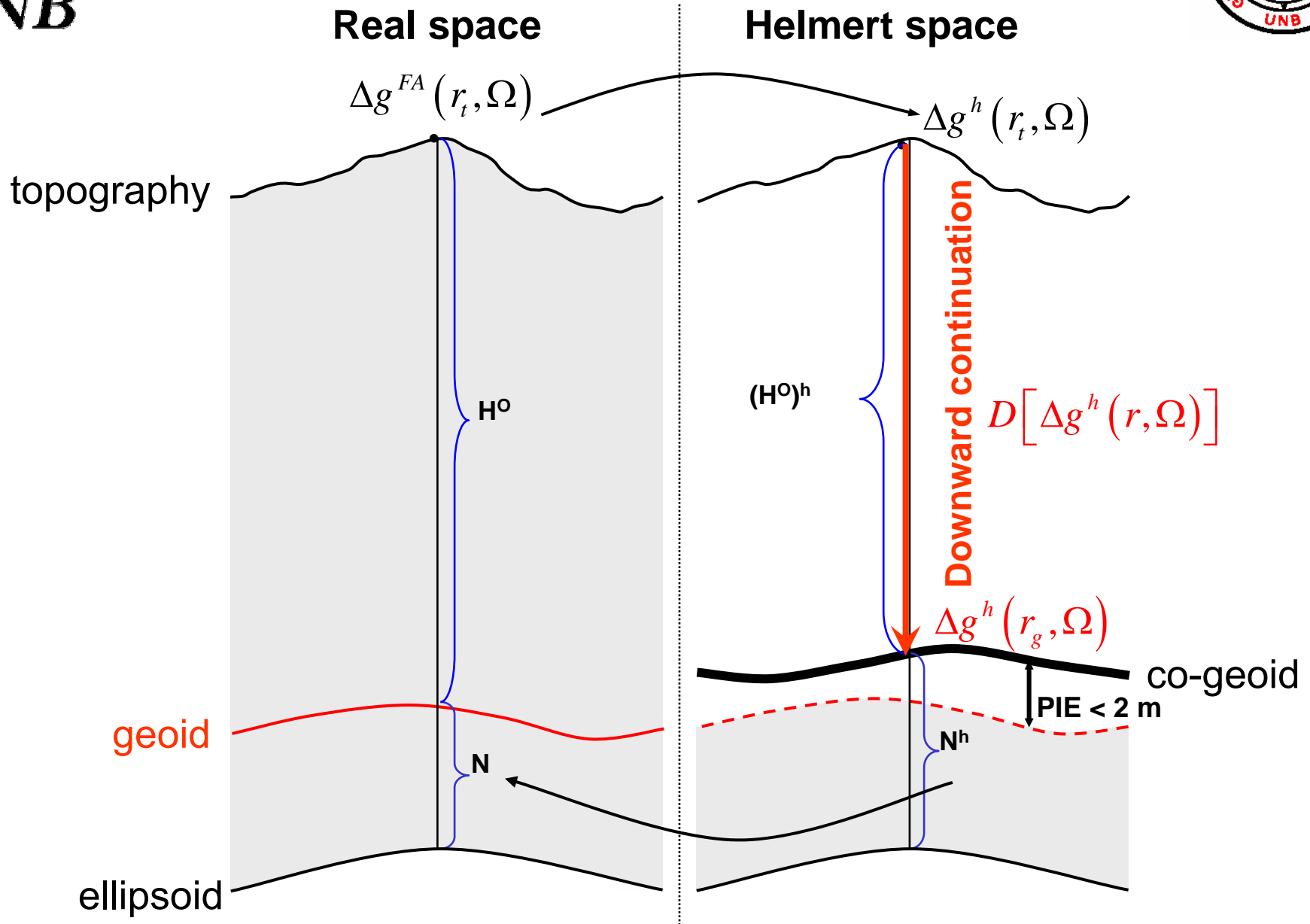
$$A(r_t, \Omega) = 4\pi G\rho H + O(H/R)$$

$$A(r_g, \Omega) = 0$$

$$V^t(r_g, \Omega) = 4\pi G\rho RH + O(H/R)$$

Rigorous far-zone contribution

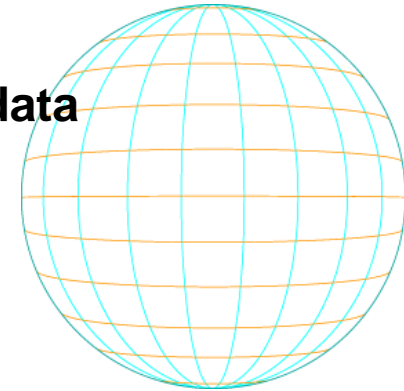
Meaningful condensation models: preserves mass or mass-centre



# Stokes's solution

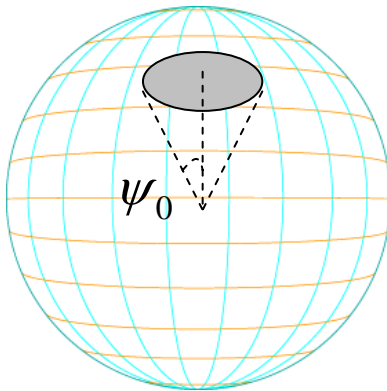
The original formula requires global coverage of gravity data

$$N(\Omega) = \frac{R}{4\pi\gamma_0(\phi)} \iint_{\Omega \in \Omega_0} \Delta g^h(R, \Omega') S(\psi(\Omega, \Omega')) d\Omega' + N_I(\Omega)$$

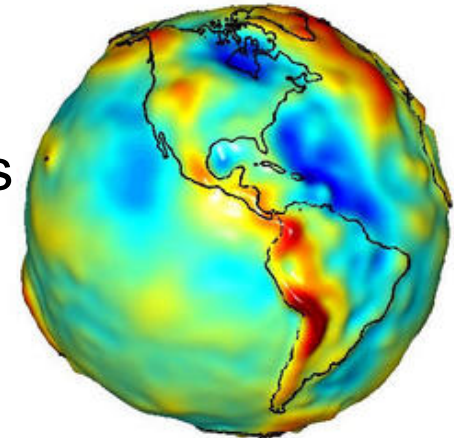


Combination of a high-degree reference field and GGM

$$N(\Omega) = \frac{R}{4\pi\gamma_0(\phi)} \iint_{\Omega \in \Omega_{\psi_0}} \left[ \Delta g^h(R, \Omega') - \sum_{n=2}^M \Delta g_n^h(\Omega) \right] S^L(\psi(\Omega, \Omega')) d\Omega' + N_I(\Omega) + \frac{R}{2\gamma} \sum_{n=2}^M \frac{2}{n-1} \Delta g_n^h(\Omega)$$



Near and far-zone contributions



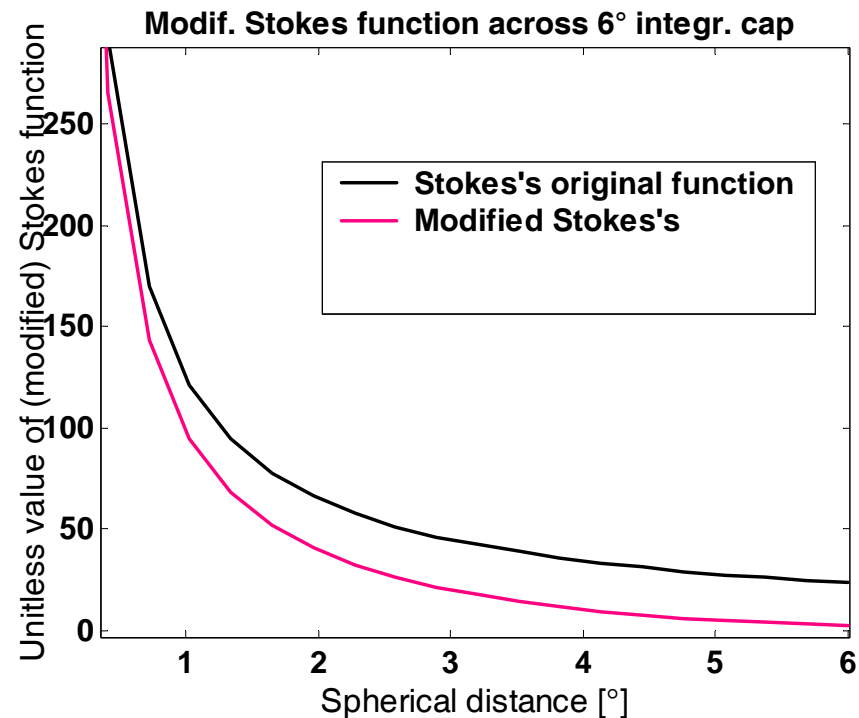
$$S^L(\psi(\Omega, \Omega')) = \sum_{n=2}^{\infty} \frac{2n+1}{n-1} P_n(\cos \psi) - \sum_{n=2}^L \frac{2n+1}{2} s_n P_n(\cos \psi)$$

$s_n \approx \frac{2}{n-1}$

$$S(\psi(\Omega, \Omega')) = \sum_{n=2}^{\infty} \frac{2n+1}{n-1} P_n(\cos \psi)$$

## Truncation bias

$$\frac{R}{2\gamma} \sum_{M+1}^{120} Q_n^L \Delta g_n(\Omega)$$



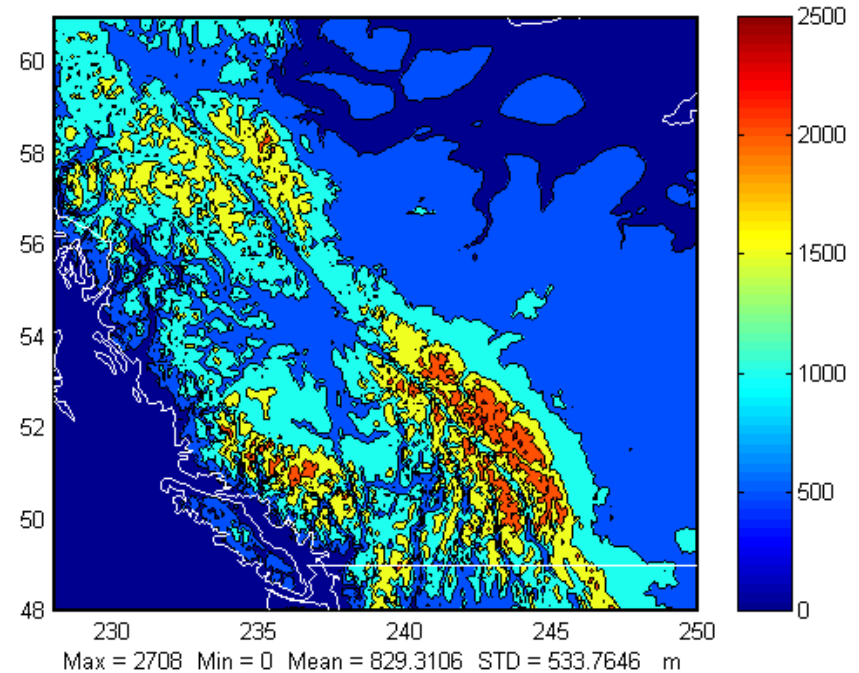


## Comparison with GPS-leveling data:

Low-land - STD < 5...10 cm,  
Mountains - STD 12-13 cm

## Data:

Gravity data resolution 5'x5',  
DTMs: 3"x3"; 30"x30"; 5'x5'  
Mean Helmert anomalies



## Testing:

Synthetic data: AUS-SEGM

2'x2' gravity data resolution

New anomaly type – NoTopography anomaly (smoother!)

Laterally varying density

CHAMP & GRACE-based geopotential models



# Conclusions



*If time permits....*

Textbook in preparation

P. Vaniček, Z. Martinec, L. Sjöberg

*Tentative title:* “The Geoid”

Springer Verlag 2006 (?)